ECONOMIC DIFFERENCE SCHEMES FOR A HYPERBOLIC SYSTEM OF EQUATIONS WITH COMPOUND DERIVATIVES AND THEIR APPLICATION TO EQUATIONS IN THE THEORY OF ELASTICITY*

A.A. SAMARSKII Moscow

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1. IN this section we consider first of all additive difference schemes (see [1] -[5]) for a system of second order hyperbolic equations which contain compound derivatives. There are schemes with variable directions which are absolutely stable and convergent at least with a speed

$$O(|h|^2+\tau)$$
, where $|h|^2=\sum_{\alpha=1}^p h_{\alpha}^2$, h_{α} is the step in the variable a_{α}

and p is the number of dimensions. The numerical algorithm consists of the conversion of a three-point triangular operator, which reduces to the successive application of known formulae. The number of operations to determine a solution for a new time layer is proportional to the number of nodes of the space network and is a quantity of the same order as the number of operations for a purely explicit scheme. Thus the schemes put forward below are economic.

To construct economic schemes for an equation of the form

$$\frac{\partial^2 u}{\partial t^2} + \sum_{\alpha, \beta=1}^{\mathbf{p}} A_{\alpha\beta}(t)u = f$$

we use the common property of the operator

$$A=\sum_{\alpha,\,\beta=1}^{p}A_{\alpha\beta},$$

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i.e. its representation in the form of the sum of operators $A_{\alpha\beta}$ of simpler structure.

Additive economic schemes for a general second order hyperbolic system are then used to describe economic schemes for a system of equations in the theory of elasticity in the case of two or three space variables (p = 2, p = 3).

A resolving scheme is also constructed for equations of elasticity, which is absolutely stable and converges with a speed $O(|h|^2 + \tau^2)$. With regard to economy the resolving scheme is comparable with additive schemes, but for convergence it requires more smoothness for the solution of the differential equation.

In Section 7 an iteration scheme with alternating directions for the solution of a difference problem, corresponding to a stationary problem in the theory of elasticity, is considered.

The convergence of this scheme is proved for p=2, 3, and it is shown that the number of iterations $v=O(h^{-2(p-1)/p}\ln{(1/\epsilon)})$, where ϵ is the required accuracy.

Economic schemes of another type are considered in the two-dimensional case (p = 2) for a dynamic problem of the theory of elasticity in [6] and for a static problem in the theory of elasticity in [7].

2. Let $\overline{G} = G + \Gamma = \{0 \leqslant x_{\alpha} \leqslant l_{\alpha}, \alpha = 1, ..., p\}$ be a p-dimensional parallelepiped with boundary Γ . In the cylinder $\overline{Q}_T = \overline{G} \times [0 \leqslant t \leqslant T]$ there is a solution of the problem

$$\frac{\partial^{2}\mathbf{u}}{\partial t^{2}} = \sum_{\alpha, \beta=1}^{p} L_{\alpha\beta}\mathbf{u} + \mathbf{f}(x, t), \qquad L_{\alpha\beta}\mathbf{u} = \frac{\partial}{\partial x_{\alpha}} \left(k_{\alpha\beta}(x, t) \frac{\partial \mathbf{u}}{\partial x_{\beta}} \right), \qquad (1)$$

$$\mathbf{u}|_{\Gamma} = \mathbf{v}(x, t), \qquad 0 \leqslant t \leqslant T, \quad \mathbf{u}(x, 0) = \mathbf{v}_{0}(x),$$

$$\frac{\partial \mathbf{u}}{\partial t}(x, 0) = \mathbf{v}_{1}(x), \qquad x \in \overline{G}.$$

Here $x=(x_1,\ldots,x_p)$; $\mathbf{u}=\mathbf{u}(x,t)=(u^1,\ldots,u^s,\ldots,u^n)$, $\mathbf{f},\mathbf{v},\mathbf{v}_0,\mathbf{v}_1$ are vectors of dimensionaly n, and $k=(k_{\alpha\beta})=(k_{\alpha\beta}{}^{sm})$, s, $m=1,\ldots,n$, is a partitioned $p\times p$ matrix with $n\times n$ submatrices, which satisfies the symmetry condition

$$k_{\alpha\beta}^{sn}(x,t) = k_{\beta\alpha}^{ms}(x,t) \tag{3}$$

and the condition of positive definiteness

$$\sum_{s, m=1}^{n} \sum_{\alpha, \beta=1}^{p} k_{\alpha\beta}^{sm}(x, t) \xi_{\beta}^{m} \xi_{\alpha}^{n} \geqslant c_{1} \sum_{s=1}^{n} \sum_{\alpha=1}^{p} (\xi_{\alpha}^{s})^{2}, \quad (x, t) \in \overline{Q}_{T}, \quad (4)$$

where $\xi_{\alpha} = (\xi_{\alpha}^{1}, \ldots, \xi_{\alpha}^{s}, \ldots, \xi_{\alpha}^{n}) \not\equiv 0$ is an arbitrary real vector and c_{1} is a positive constant.

We shall assume that the problem (1) - (2) has a unique solution $u=u(x,\ t)$, which is continuous in Q_T and differentiable as many times as necessary for this work. To infer the *a priori* evaluations it is assumed that $k_{\alpha\beta}(x,t)$ satisfies a Lipschitz condition with respect to t and $x_{\alpha'}$, $\alpha'=1,\ldots,p$.

The system of equations in the theory of elasticity

$$\frac{\partial^2 \mathbf{u}}{\partial t^2} = \mu \Delta \mathbf{u} + (\lambda + \mu) \operatorname{grad} \operatorname{div} \mathbf{u} + \mathbf{f}(x, t), \tag{5}$$

where $\Delta \mathbf{u} = \sum_{\alpha=1}^p \partial^2 \mathbf{u} / \partial x_{\alpha}^2$ is the Laplace operator, $\mathbf{u} = (u^1, \ldots, u^p)$,

 $\lambda = {\rm const} > 0$ and $\mu = {\rm const} > {\rm are\ Lam\'e's\ coefficients}$, is obviously a particular case of the system of equations (1) with n=p and

$$k_{\alpha\beta}^{sm} = \mu \delta_{\alpha\beta} \delta_{sm} + (\lambda + \mu) \delta_{\alpha s} \delta_{\beta m}, \qquad \delta_{ij} = \begin{cases} 1, & i = j, \\ 0, & i \neq j, \end{cases}$$
 (6)

where δ_{ij} is the Kronecker delta. Condition (3) is satisfied automatically. We shall show that condition (4) is also satisfied if $c_1 = \mu$. In fact

$$\sum_{s, m=1}^{p} \sum_{\alpha, \beta=1}^{p} k_{\alpha\beta}^{m} \xi_{\alpha}^{s} \xi_{\beta}^{m} = \mu \sum_{\alpha, s=1}^{p} (\xi_{\alpha}^{s})^{2} + (\lambda + \mu) \sum_{\alpha, s=1}^{p} \xi_{\alpha}^{\alpha} \xi_{s}^{s} =$$

$$=\mu\sum_{\alpha, s=1}^{p}(\xi_{\alpha}^{s})^{2}+(\lambda+\mu)\left(\sum_{\alpha=1}^{p}\xi_{\alpha}^{\alpha}\right)^{2}\geqslant\mu\sum_{\alpha, s=1}^{p}(\xi_{\alpha}^{s})^{2}.$$

3. Let us introduce the difference networks $\omega_{\tau} = \{t_j = j\tau \in [0, T], j = 0, 1, \ldots\}$ and $\overline{\omega}_h = \{x_i = (i_1h_1, \ldots, i_{\alpha}h_{\alpha}, \ldots, i_{p}h_p) \in \overline{G} = G + \Gamma;$ $i_{\alpha} = 0, 1, \ldots, N_{\alpha}, h_{\alpha} = l_{\alpha}/N_{\alpha}, \alpha = 1, 2, \ldots, p\}$ with steps τ for the

variable t and h_{α} for the variable x_{α} , $\alpha = 1, \ldots, p$: let $\gamma_h = \{x_i \in \Gamma\}$ be the boundary of the network $\overline{\omega}_h$, $\overline{\omega}_h \setminus \gamma_h = \{x_i \in G\}$ be the set of inner

nodes, $|h|^2 = \sum_{\alpha=1}^p h_{\alpha}^2$. Following [1], we shall introduce the notation $y = y(x_i, t_{i+1}) = y^{j+1}$, $y = y^j$,

$$x_i^{(\pm 1_{\alpha})} = (i_1 h_1, \dots, i_{\alpha-1} h_{\alpha-1}, (i_{\alpha} \pm 1) h_{\alpha}, i_{\alpha+1} h_{\alpha+1}, \dots, i_p h_p),$$

$$\mathbf{y}^{(\pm 1_{\alpha})} = \mathbf{y} \left(x_i^{(\pm 1_{\alpha})}, t_{j+1} \right), \quad \mathbf{y}_{\bar{x}_{\alpha}} = \left(\mathbf{y} - \mathbf{y}^{(-1_{\alpha})} \right) / h_{\alpha}, \quad \mathbf{y}_{x_{\alpha}} = \left(\mathbf{y}^{(+1_{\alpha})} - \mathbf{y} \right) / h_{\alpha}.$$

We shall replace the operator

$$L_{\alpha\beta}\mathbf{u} = \frac{\partial}{\partial x_{\alpha}} \left(k_{\alpha\beta} (x, t) \frac{\partial \mathbf{u}}{\partial x_{\beta}} \right) \tag{7}$$

in the difference network ω_h by the same scheme of the second order of approximation as in [2], assuming that

$$\Lambda_{\alpha\beta}\mathbf{y} = \frac{1}{2} \left[\left(a_{\alpha\beta} \mathbf{y}_{\overline{x}_{\beta}} \right)_{x_{\alpha}} + \left(a_{\alpha\beta}^{(+1\beta)} \mathbf{y}_{x_{\beta}} \right)_{\overline{x}_{\alpha}} \right] \text{ for } \beta \neq \alpha,
\Lambda_{\alpha\alpha}\mathbf{y} = \left(a_{\alpha\alpha} \mathbf{y}_{\overline{x}_{\alpha}} \right)_{x_{\alpha}},$$
(8)

where $(a_{\alpha\beta})$ is a matrix-functional of the matrix $(k_{\alpha\beta})$ with pattern $\{-1\leqslant s_{\beta}\leqslant 0,\ \beta=1,\ldots,p\}$. The coefficients $a_{\alpha\beta}=(a_{\alpha\beta}^{sm})$ satisfy the conditions

$$a_{\alpha\beta}^{sm} = a_{\beta\alpha}^{ms}, \tag{9}$$

for sufficiently small $|h| \leqslant h_0$, the condition

$$\sum_{s,m=1}^{n} \sum_{\alpha,\beta=1}^{p} a_{\alpha\beta}^{sm} \xi_{\beta}^{m} \xi_{\alpha}^{s} \geqslant c_{1}' \sum_{\alpha=1}^{p} \sum_{s=1}^{n} (\xi_{\alpha}^{s})^{2}, \quad (x,t) \in \overline{\omega}_{h} \times \omega_{\tau}, \quad (10)$$

where $c_1 \leq c_1$ is a positive constant which does not depend on the network, and the condition obtained from (10) after replacing $a_{\alpha\beta}^{sm}$ by the coefficinet $(a_{\beta}^{sm})^{(1\beta)}$ at the point $x_i^{(+1)\beta}$.

In the case of constant coefficients, $k_{\alpha\beta}=$ const., obviously $a_{\alpha\beta}=k_{\alpha\beta}$ and instead of (8) we obtain

$$\Lambda_{\alpha\beta} \mathbf{y} = \frac{1}{2} k_{\alpha\beta} (\mathbf{y}_{\overline{\mathbf{x}}_{\beta} \mathbf{x}_{\alpha}} + \mathbf{y}_{\overline{\mathbf{x}}_{\alpha} \mathbf{x}_{\beta}}), \qquad \Lambda_{\alpha\alpha} \mathbf{y} = k_{\alpha\alpha} \mathbf{y}_{\mathbf{x}_{\alpha} \mathbf{x}_{\alpha}}; \tag{11}$$

condition (10) is satisfied on any network.

Note. For $\Lambda_{\alpha\beta}y$ instead of (8) we can consider other representations also, e.g.

$$\Lambda_{\alpha\beta}y = \frac{1}{2} \left[\left(a_{\alpha\beta} \mathbf{y}_{\bar{\alpha}\beta} \right)_{\bar{x}_{\alpha}} + \left(a_{\alpha\beta}^{(+1)\beta} \mathbf{y}_{\mathbf{x}\beta} \right)_{\mathbf{x}_{\alpha}} \right],$$

$$\Lambda_{\alpha\beta}y = \frac{1}{2} \left[\left(a_{\alpha\beta} \mathbf{y}_{\bar{x}\beta} \right)_{\hat{x}_{\alpha}}^{*} + \left(a_{\alpha\beta}^{(+1)\beta} \mathbf{y}_{\mathbf{x}\beta} \right)_{\hat{x}_{\alpha}}^{*} \right],$$
(12)

where $\mathbf{v}_{\hat{x}_{\alpha}} = \frac{1}{2} (\mathbf{v}_{\bar{x}_{\alpha}} + \mathbf{v}_{x_{\alpha}})$ is the central difference derivative. In all cases the condition (10) will be satisfied for a sufficiently small step $|h| \leqslant h_0$ and all the subsequent conclusions retain their validity.

4. Let us introduce the "triangular" operators L^- and L^+ . For this we write the symmetrical matrix $k_{\alpha\alpha}=(k_{\alpha\alpha}^{sm})$ in the form of the sum of two triangular matrices $k_{\alpha\alpha}=k_{\alpha\alpha}^-+k_{\alpha\alpha}^+, k_{\alpha\alpha}^-=(k_{\alpha\alpha}^{-sm}), k_{\alpha\alpha}^+=(k_{\alpha\alpha}^{+sm}),$ assuming $k_{\alpha\alpha}^{-ss}=k_{\alpha\alpha}^{+ss}=\frac{1}{2}k_{\alpha\alpha}^{ss}, k_{\alpha\alpha}^{-sm}=k_{\alpha\alpha}^{sm}, k_{\alpha\alpha}^{+sm}=0$ if $m < s, k_{\alpha\alpha}^{+sm}=k_{\alpha\alpha}^{sm}, k_{\alpha\alpha}^{-sm}=0$ with m > s and any $\alpha = 1, \ldots, p$. The matrix $k_{\alpha\alpha}^+$ is a diagonal $p \times p$ matrix with submatrices which are triangular $n \times n$ matrices, conjugate to each other

$$k_{\alpha\alpha}^{-sm} = k_{\alpha\alpha}^{+ms} \qquad (a_{\alpha\alpha}^{-sm} = a_{\alpha\alpha}^{+ms}). \tag{13}$$

In accordance with the representation $k_{lphalpha}=k_{lphalpha}^-+k_{lphalpha}^+$ we obtain

$$L_{\alpha\alpha} = L_{\alpha\alpha}^- + L_{\alpha\alpha}^+, \quad \Lambda_{\alpha\alpha} = \Lambda_{\alpha\alpha}^- + \Lambda_{\alpha\alpha}^+,$$

where

$$L_{lphalpha}^{}\mathbf{u}=rac{\partial}{\partial x_{lpha}}\Big(k_{lphalpha}^{}rac{\partial\mathbf{u}}{\partial x_{lpha}}\Big)$$

etc. In view of condition (13) the operators $\Lambda_{\alpha\alpha}^-$ and $\Lambda_{\alpha\alpha}^+$ are conjugate to each other on the network ω_h in the sense of the scalar product

$$(\mathbf{y}, \mathbf{v}) = \sum_{\omega_h} \mathbf{y}(x_i) \mathbf{v}(x_i) H, \quad H = h_1 \dots h_p,$$

i.e.

$$(\Lambda_{\alpha\alpha}^-\mathbf{y}, \mathbf{v}) = (\Lambda_{\alpha\alpha}^+\mathbf{v}, \mathbf{y}), \quad a = 1, \dots p,$$
 (14)

where y and v are arbitrary network functions, vanishing on the boundary γ_h of the network ω_h .

We shall put the operator

$$L = \sum_{\alpha, \beta=1}^{\mathbf{p}} L_{\alpha\beta}$$

in the form of the sum of two triangular operators

$$L = L^{-} + L^{+}, \qquad L^{\mp} = \sum_{\alpha, \beta=1}^{p} L_{\alpha\beta}^{\mp} = \sum_{\alpha=1}^{p} L_{\alpha}^{\mp}, \qquad L_{\alpha}^{\mp} = \sum_{\beta=1}^{p} L_{\alpha\beta}^{\mp},$$

$$L_{\alpha\beta}^{\mp} = L_{\alpha\alpha}^{\mp} \text{ if } \beta = \alpha, \qquad L_{\alpha\beta} = L_{\alpha\beta},$$

$$L_{\alpha\beta}^{+} = 0 \text{ if } \beta \leqslant \alpha, \qquad L_{\alpha\beta}^{+} = L_{\alpha\beta},$$

$$L_{\alpha\beta} = 0 \text{ if } \beta > \alpha, \qquad L_{\alpha}^{-} = L_{\alpha\alpha} + \sum_{\beta=1}^{p} L_{\alpha\beta}, \qquad L_{\alpha} = L_{\alpha\alpha} + \sum_{\beta=\alpha+1}^{p} L_{\alpha\beta}.$$

$$(15)$$

By virtue of the principle of additivity of [1 - 5] the solution of the system of equations

$$\frac{\partial^2 \mathbf{u}}{\partial t^2} = \sum_{\alpha=1}^p [(L_{\alpha}^- + L_{\alpha}^+)\mathbf{u} + \mathbf{f}_{\alpha}], \qquad \sum_{\alpha=1}^p \mathbf{f}_{\alpha} = \mathbf{f}$$

reduces to the successive solution, on an $\omega_h \times \omega_T$ network with step τ/p , of the simpler equations

$$\frac{1}{p}\frac{\partial^2 \mathbf{u}}{\partial t^2} = L_{\alpha}^{-}\mathbf{u} + L_{\alpha}^{+}\mathbf{u} + \mathbf{f}_{\alpha}. \tag{16}$$

The case where $L_{\alpha\beta} = \delta_{\alpha\beta}L_{\alpha\alpha}$, i.e. compound derivatives are absent, is considered in [1].

We now introduce the values $y^{j+\alpha/p} = y_{(\alpha)}$, intermediate between $y^j = y$ and $y^{j+1} = y$ assuming that $y^{(j-1)+\alpha/p} = y_{(\alpha)}$, and use for the determination of $y_{(\alpha)}$ a difference scheme which approximates equation (16) with number α . By analogy with [1], for the approximation of $\partial^2 u / \partial t^2$ we use the (p+1)-th time layer

$$\frac{1}{p} \frac{\partial^2 \mathbf{u}}{\partial t^2} \sim \sigma_p \mathbf{u}_{\bar{t}_\alpha \bar{t}_\alpha}, \qquad \alpha = 1, \ldots, p; \tag{17}$$

$$\mathbf{u}_{\overline{t}_{\alpha}\overline{t}_{\alpha}} = \begin{cases} \left(\mathbf{u}_{(\alpha)} - 2\mathbf{u}_{(\alpha-1)} + \check{\mathbf{u}}_{(\alpha)}\right)/\tau^{2}, & \sigma_{p} = 2, & \mathbf{u}_{(0)} = \check{\mathbf{u}}_{(2)}, & p = 2 \\ \\ \left(\mathbf{u}_{(\alpha)} - \mathbf{u}_{(\alpha-1)} - \mathbf{u}_{(\alpha-2)} + \check{\mathbf{u}}_{(\alpha)}\right)/\tau^{2}, & \sigma_{p} = \frac{3}{2}, & \mathbf{u}_{(0)} = \check{\mathbf{u}}_{(3)}, \\ \\ \mathbf{u}_{(-1)} = \check{\mathbf{u}}_{(2)}, & p = 3. \end{cases}$$

The additive scheme of alternating directions for the problem (1) - (2) will have the form

$$\sigma_{p}\mathbf{y}_{\bar{t}_{\alpha}\bar{t}_{\alpha}} = \sum_{\beta=1}^{\alpha} \Lambda_{\alpha\beta}^{-}\mathbf{y}_{(\beta)} + \sum_{\beta=\alpha}^{p} \Lambda_{\alpha\beta}^{+} \dot{\mathbf{y}}_{(\beta)} + \varphi_{\alpha}, \quad \alpha = 1, \ldots, p, \qquad (x, t) \in \omega_{h} \times \omega_{t},$$
(18)

$$y_{\alpha} = y(x, t_{j}^{*}) \text{ with } x_{\alpha} = 0, l_{\alpha}, \qquad \alpha = 1, \ldots, p; \qquad y(x, 0) = y_{0}(x), \quad (19)$$

where $\phi_{\alpha} = \phi_{\alpha}(x, t_{j}^{*})$ is a second order approximation on the network ω_{h} of the function $f_{\alpha}(x, t), t_{j}^{*} \in [t_{j}, t_{j+1}], \text{ e.g. } t_{j}^{*} = t_{j+1/2} = t_{j} + 0.5 \tau.$ The coefficients $\alpha_{\alpha\beta} = \alpha_{\alpha\beta}(x, t_{(\alpha)}^{*})$ are taken at the middle moment $t_{(\alpha)}^{*} = t_{j} + \frac{\alpha}{2}\tau - 0.5\tau$ (cf. [1]).

The second initial condition can be approximated by analogy with [1], or more simply by assuming

$$\mathbf{y}^{\alpha/p} = \mathbf{v}_1(x) + \frac{\alpha \tau}{\rho} \tilde{\mathbf{v}}_1(x), \quad \alpha = 1, ..., p-1, \quad p = 2, 3.$$
 (20)

Such a condition is sufficient for an accuracy $O(\tau + |h|^2)$. Let us rewrite (18) in the form

$$\left(E - \frac{\tau^2}{\sigma_p} \Lambda_{\alpha\alpha}^{-}\right) \mathbf{y}_{(\alpha)} = R_{\alpha} (\mathbf{y}) + \frac{\tau^2}{\sigma_p} \sum_{\beta=1}^{\alpha-1} \Lambda_{\alpha\beta} \mathbf{y}_{(\beta)} + \mathbf{F}_{\alpha} = \mathbf{\Phi}_{\alpha}, \tag{21}$$

where

$$\mathbf{F}_{\alpha} = \frac{\tau^{2}}{\sigma_{p}} \left[\sum_{\beta=\alpha}^{p} \Lambda_{\alpha\beta}^{+} \dot{\mathbf{y}}_{(\beta)}^{\cdot} + \varphi_{\alpha} \right],$$

 $R_{\alpha}(\mathbf{y}) = 2\mathbf{y}_{(\alpha-1)} - \mathbf{y}_{(\alpha)}$ if p = 2, $R_{\alpha}(\mathbf{y}) = \mathbf{y}_{(\alpha-1)} + \mathbf{y}_{(\alpha-2)} - \mathbf{y}_{(\alpha)}$ if p = 3 and E is the operator of identity.

Here it is obvious that to determined $y_{(\alpha)}$ (all $y_{(\beta)}$ for $\beta \leq \alpha$ and all $y_{(\beta)}$ for $\beta = 1, \ldots, p$ are already known) we obtain a system of three-

point equations with a triangular matrix for their coefficients. Its solution reduces to an inversion of the operator $E = (\tau^2/\sigma_p)\Lambda_{\alpha\alpha}$, which is attained by an n-fold use of the ordinary formulae for each chain $y_{(\alpha)}$ (see [1]) for fixed $\alpha = 1, \ldots, p$. To realize the algorithm (21) we must bear in mind the values of $y_{(\alpha)}$ on p layers.

Let us write the equation for the s-th components $y_{(\alpha)}^s$ of the vector $\mathbf{y}_{(\alpha)}$

$$y_{(\alpha)}^{s} - \frac{\tau^{s}}{\sigma_{p}} (a_{\alpha\alpha}^{-ss} y_{\widetilde{x}_{\alpha}}^{s})_{x_{\alpha}} = \Phi_{(\alpha)}^{s} + \frac{\tau^{s}}{\sigma_{p}} \sum_{m=1}^{s-1} (a_{\alpha\alpha}^{-sm} y_{\widetilde{x}_{\alpha}}^{m})_{x_{\alpha}}. \tag{22}$$

Here $\Phi_{\alpha}{}^s$ is known, since the calculation takes place in the direction from α to $\alpha+1$, $\alpha=1,\ldots,p$; the second term is also known if we determine successively $y_{(\alpha)}{}^i,\ldots,y_{(\alpha)}{}^s,\ldots,y_{(\alpha)}{}^p$, i.e. carry out the calculation from s to s+1. Hence it is obvious that we can find $y_{(\alpha)}{}^s$ by solving the first boundary value problems for the three-point equations on segments parallel to the axis Ox_{α} .

If we interchange the roles of L_{α} and L_{α} , respectively,

$$\Lambda_{\alpha}^- = \Lambda_{\alpha\alpha}^- + \sum_{\beta=1}^{\alpha-1} \Lambda_{\alpha\beta} = \sum_{\beta=1}^{\alpha} \Lambda_{\alpha\beta}^- \text{ and } \Lambda_{\alpha}^+ = \Lambda_{\alpha\alpha}^+ + \sum_{\beta=\alpha+1}^p \Lambda_{\alpha\beta} = \sum_{\beta=\alpha}^p \Lambda_{\alpha\beta}^+,$$

we obtain a second additive scheme

$$\sigma_p \mathbf{y}_{\vec{t}_{\alpha}\vec{t}_{\alpha}} = \sum_{\beta=\alpha}^{p} \Lambda_{\alpha\beta}^{+} \mathbf{y}_{(\beta)} + \sum_{\beta=1}^{\alpha} \Lambda_{\alpha\beta}^{+} \dot{\mathbf{y}}_{(\beta)} + \varphi_{\alpha}$$
 (23)

with the same initial and boundary conditions as in the first scheme. Here to determine y_{α} we must invert the triangular three-point operator $E = (\tau^2/\sigma_p)\Lambda_{\alpha\alpha}^+$. Here the calculation proceeds from $\alpha + 1$ to α and from s + 1 to s.

The alternation of these two schemes gives a third scheme. Introducing the intermediate value $y^{j+\alpha/2p}$, $\alpha=1,\ldots,2p-1$, we obtain (cf. [2, 5])

$$\sigma_{p}\mathbf{y}_{\bar{t}_{\alpha}\bar{t}_{\alpha}} = \sum_{\beta=1}^{\alpha} \Lambda_{\alpha\beta}^{-}\mathbf{y}_{(\beta)} + \sum_{\beta=\alpha}^{p} \Lambda_{\alpha\beta}^{+}\dot{\mathbf{y}}_{(\beta)} + \varphi_{\alpha}, \qquad \alpha = 1, \ldots, p;$$
(24)

$$\sigma_{p}\mathbf{y}_{\bar{t}_{\alpha'}\bar{t}_{\alpha'}} = \sum_{\beta=\alpha}^{p} \Lambda_{\alpha\beta}^{+}\mathbf{y}_{(\beta')} + \sum_{\beta=1}^{\alpha} \Lambda_{\alpha\beta'}^{-}\mathbf{y}_{(\beta')} + \varphi_{\alpha'},$$

$$\alpha' = 2p + 1 - \alpha, \quad \beta' = 2p + 1 - \beta;$$

where
$$\alpha' = p + 1, ..., 2p$$
, $\varphi_{\alpha'} = \varphi_{\alpha}$, $\alpha = p, p - 1, ..., 2, 1$.

5. Schemes (18) and (23) are stable for sufficiently small $|h| \leq h_0$, ensuring that the requirement (10) of the positive definiteness of the matrices $(a_{\alpha\beta})$ and $(a_{\alpha\beta}^{(+1\beta)})$ for any τ is satisfied, and converge at least with a speed $O(|h|^2 + \tau)$. The proof of these statements is carried out by analogy with [1, 4] by the method of energy inequalities.

Here the basic part is played by an identity of the form

$$\sum_{\alpha=1}^{p} \{ (a_{\alpha\beta}^{-} \xi_{\beta}, \xi_{\alpha} - \xi_{\alpha}) + (a_{\alpha\beta}^{+} \xi_{\beta}, \xi_{\alpha} - \xi_{\alpha}) = J - \check{J} (1 + O(\tau)) + R,$$

where

$$J = \sum_{\alpha=1}^{p} \sum_{\beta=1}^{\alpha} (a_{\alpha\beta}^{-} \xi_{\beta}, \xi_{\alpha}) = \sum_{\alpha=1}^{p} \sum_{\beta=\alpha}^{p} (a_{\alpha\beta}^{+} \xi_{\beta}, \xi_{\alpha}),$$

$$R = \sum_{\alpha=1}^{p} \left\{ \sum_{\beta=1}^{\alpha} (a_{\alpha\beta}^{+} \xi_{\beta}, \xi_{\alpha}) - (a_{\alpha\beta}^{-} \xi_{\beta}, \xi_{\alpha}) \right\}.$$

Using (3) it is not difficult to see that R = 0.

Note. In the case of constant coefficients. $k_{\alpha\beta}=$ const., the given values of the accuracy of the schemes considered, (18) and (23), are valid for any h_{α} and τ .

6. We now turn to the equations of elasticity (5). In this case as we have seen in Section 2, n = p, and $k_{\alpha\beta}^{sm} = \text{const.}$,

$$k_{\alpha\beta}^{sm} = \mu \delta_{\alpha\beta} \delta_{sm} + (\lambda + \mu) \delta_{\alpha s} \delta_{\beta m}, \qquad \delta_{sm} = \begin{cases} 1, & s = m, \\ 0, & s \neq m. \end{cases}$$
 (25)

For the equations of elasticity (5) we can use any of the schemes considered in Section 4, bearing in mind that $a_{\alpha\beta}=k_{\alpha\beta}$, where $k_{\alpha\beta}$ is determined by formula (25).

We shall write in more detail the difference equations (18) in the

case where p = 2; here $y = (y^{(1)}, y^{(2)})$ and

$$2y_{\overline{t_1}\overline{t_1}}^{(1)} = \frac{1}{2}(\lambda + 2\mu) \left(y_{(1)}^{(1)} \bar{x_i} x_i + \check{y}_{(1)}^{(1)} \bar{x_i} x_i\right) + \frac{1}{2}(\lambda + \mu) \left(\check{y}_{(2)}^{(2)} \bar{x_s} x_i + \check{y}_{(2)}^{(2)} x_s \bar{x_i}\right) + \varphi_{(1)}^{(1)}$$

$$2y_{\overline{t_1}\overline{t_1}}^{(2)} = \frac{1}{2} \mu \left(y_{(1)}^{(2)} \frac{1}{x_1 x_1} + \tilde{y}_{(1)}^{(2)} \frac{1}{x_1 x_1} \right) + \varphi_{(1)}^{(2)},$$

$$2y_{\overline{t_2t_2}}^{(1)} = \frac{1}{2} \mu \left(y_{(2)}^{(1)} \bar{x_2} x_2 + y_{(2)}^{(1)} \bar{x_2} x_2 \right) + \phi_{(2)}^{(1)},$$

$$2y_{\overline{t_1t_2}}^{(2)} = \frac{1}{2} (\lambda + 2\mu) (y_{(2)}^{(2)} \frac{1}{x_2x_2} + \tilde{y}_{(2)}^{(2)} \frac{1}{x_2x_2}) + \frac{1}{2} (\lambda + \mu) (y_{(1)}^{(1)} \frac{1}{x_2x_1} + y_{(1)}^{(1)} \frac{1}{x_1x_2}) + \Phi_{(2)}^{(2)}.$$

We must remember that here the upper index means the number of the component and the lower one the number of the vector $(y_{(1)}^{(1)} = (y^{(1)})^{j+1/2}$, $y_{(2)}^{(1)} = (y^{(1)})^{j+1} = y^{(1)}$, $y_{(2)}^{(2)} = (y^{(2)})^{j+1} = y^{(2)}$ etc.). Using (18) and (25) it is not difficult to write down the scheme for p = 3.

7. For the equations of elasticity we can also construct a series of resolving schemes which are absolutely stable, economic and convergent with speed $O(\tau + |h|^2)$ or $O(\tau^2 + |h|^2)$.

Let us first consider the two multidimensional schemes

$$\mathbf{y}_{\bar{t}\bar{t}} = \Lambda^{-}\mathbf{y} + \Lambda^{+}\check{\mathbf{y}} + \check{\mathbf{\varphi}} \qquad (\mathbf{y}_{\bar{t}\bar{t}} = (\mathbf{y}^{j+1} - 2\mathbf{y}^{j} + \mathbf{y}^{j-1})/\tau^{2}); \tag{I}$$

$$\mathbf{y}_{\bar{t}\bar{t}} = \mathbf{\Lambda}^{+}\mathbf{y} + \mathbf{\Lambda}^{-}\mathbf{\tilde{y}} + \mathbf{\tilde{\phi}}, \tag{I}^{+}$$

where Λ^- and Λ^+ are triangular operators, which approximate the triangular differential operator L^- and L^+ , $y=y^{j+1}$, $\check y=y^{j-1}$, $\check y=y^j$.

Let $\mathring{\Lambda}_{\alpha}\mathbf{y} = \mathbf{y}_{\widetilde{x}_{\alpha}x_{\alpha}}$, and $\mathring{\Lambda}_{sk}$ denote the difference approximation of the compound derivative $\partial^{2}\mathbf{u}/\partial x_{s}\partial x_{k}$, e.g. $\mathring{\Lambda}_{sk}\mathbf{y} = \frac{1}{2}(\mathbf{y}_{\widetilde{x}_{s}x_{k}} + \mathbf{y}_{x_{s}\widetilde{x}_{k}})$ or $\mathring{\Lambda}_{sk}\mathbf{y} = \frac{1}{2}(\mathbf{y}_{\widetilde{x}_{s}x_{k}} + \mathbf{y}_{x_{s}x_{k}})$. Then the expressions for the triangular operators $\mathring{\Lambda}_{sk}\mathbf{y} = \frac{1}{2}(\mathbf{y}_{\widetilde{x}_{s}x_{k}} + \mathbf{y}_{x_{s}x_{k}})$.

ators Λ^- and Λ^+ can be written in the form (the upper index s or k is the number of the component)

$$(\Lambda^{-}\mathbf{y})^{s} = \frac{1}{2} \sum_{\alpha=1}^{p} \kappa_{s\alpha} \mathring{\Lambda}_{\alpha} y^{s} + (\lambda + \mu) \sum_{k=1}^{s-1} \mathring{\Lambda}_{sk} y^{k}, \quad \kappa_{s\alpha} = \mu + (\lambda + \mu) + \delta_{s\alpha}, \quad (26)$$

$$(\Lambda^{+}\mathbf{y})^{s} = \frac{1}{3} \sum_{\alpha=1}^{p} \varkappa_{\rho\alpha} \mathring{\Lambda}_{\alpha} y^{s} + (\lambda + \mu) \sum_{k=s+1}^{p} \mathring{\Lambda}_{sk} y^{k}.$$
 (27)

The boundary conditions on γ_h are exactly satisfied

$$\mathbf{y}|_{v_h} = \mathbf{v}(x,t),$$

and the initial conditions have the form

$$y(x, 0) = v_0(x), \quad y_{\bar{t}}(x, \tau) = v_1(x) + \tilde{v}_1(x),$$

where $\tilde{\mathbf{v}}_{\mathbf{i}}(x)$ is chosen so that the initial condition $\partial \mathbf{u}/\partial t = \mathbf{v}_{\mathbf{i}}$ is approximated with accuracy $O(\tau^2)$; for this, for instance, it is sufficient to assume that $\tilde{\mathbf{v}}_{\mathbf{i}} = -(L\mathbf{u} + \mathbf{f})|_{t=0}$.

Each of the schemes I⁻ and I⁺ is absolutely stable and has an accuracy $O(\tau + |h^2|)$. Applying these schemes alternately (e.g. scheme (26) on odd and scheme (27) on even layers), we obtain an accuracy $O(\tau^2 + |h|^2)$.

Following the principle given in (8), we shall write the generating scheme II for the scheme I

$$A^{s}y_{\bar{t}}^{s} = \check{\Phi}^{s} + F^{s}, \qquad A^{s} = \prod_{\alpha=1}^{p} A_{\alpha}^{s}, \qquad A_{\alpha}^{s} = E - 0.5\tau^{2}\kappa_{s\alpha}\mathring{\Lambda}_{\alpha},$$

$$\check{\Phi}^{s} = \check{y}_{\bar{t}}^{s} + \tau \left[(\Lambda^{+}\check{y})^{s} + 0.5 \sum_{\alpha=1}^{p} \kappa_{s\alpha}\mathring{\Lambda}_{\alpha}\check{y}^{s} + \check{\Phi}^{s} \right],$$

$$F^{s} = \tau (\lambda + \mu) \sum_{k=1}^{s-1} \mathring{\Lambda}_{sk}y^{k}.$$

$$(28)$$

Similarly the generating scheme II⁺ is written for the scheme I⁺. In this case only the formulae for $\check{\Phi^s}$ and F^s are changed

$$\check{\Phi}^{s} = \check{y}_{\tilde{t}}^{s} + \tau \left[(\Lambda^{-} \check{y})^{s} + 0.5 \sum_{\alpha=1}^{p} \varkappa_{s\alpha} \mathring{\Lambda}_{\alpha} \check{y}^{s} + \check{\Phi}^{s} \right], \qquad F^{s} = \tau (\lambda + \mu) \sum_{k=s+1}^{p} \mathring{\Lambda}_{sk} y^{k}.$$

To determine y^s on a new layer we can write some numerical algorithms for alternating directions. We give for scheme II only the algorithm which we put forward in [8] for the equation of heat conduction: this algorithm has the form

$$A_1^{s}v_{(1)}^{s} = \check{\Phi}^{s} + F^{s}, \qquad A_{\alpha}v_{(\alpha)}^{s} = v_{(\alpha-1)}^{s}, \quad \alpha = 1, 2, \dots, p, \ s = 1, \dots, p;$$

$$y^{s} = \check{y}^{s} + \tau v_{(p)}^{s}.$$

$$(30)$$

We shall take the boundary conditions with $x_{\alpha}=0$, $x_{\alpha}=l_{\alpha}$ for $v_{(\alpha)}$ in the form

$$v_{(a)}^s = A_{a+1}^s \dots A_p^s v_{\bar{l}}^s$$
 for $x_a = 0$, $x_a = l_a$, $\alpha = 1, \dots, p-1$, $v_{(p)}^s = v_{\bar{l}}^s$. (31)

The order of calculation is as follows: the components $y^{(i)}, \ldots, y^{(p)} = y^{j+1}$ are determined in turn.

For the algorithm which corresponds to scheme II⁺ the formulae (30) remain unchanged, but the order of calculation is reversed: the components $y^{(p)}, \ldots, y^{(1)}$ are determined successively.

Alternating the schemes II⁻ and II⁺ we find the solution of the problem with accuracy to within $O(|h|^2 + \tau^2)$. This evaluation is obtained by the method of [1, 2, 8].

From the formulae for $v_{(\alpha)}^s$ it is obvious that the solution y^s of the difference problem is determined by means of successive inversion of the triangular matrices (according to the formulae given in [9]). Therefore the resolving schemes are economic: to calculate the vector $|y^{j+1}|$ operations of the order $O(p^2/h^p)$ are required.

8. We turn now to a stationary problem of the theory of elasticity

$$L\mathbf{u} = \mu \Delta \mathbf{u} + (\lambda + \mu) \text{ grad div } \mathbf{u} = -\mathbf{f}(x), \quad x \in G, \quad \mathbf{u}|_{\Gamma} = \mathbf{v}(x).$$
 (32)

Its solution reduces to the solution of the difference problem for establishing the parabolic system of equations

$$\frac{\partial \mathbf{u}}{\partial t} = L\mathbf{u} + f(x), \quad \mathbf{u} \mid \mathbf{r} = \mathbf{v}(x)$$

with arbitrary initial data

$$\mathbf{u}(x,0) = \mathbf{v}_0(x).$$

For this we make use of the resolving scheme. Let the difference scheme for (32) take the form

$$\Lambda^{-}\mathbf{v} + \Lambda^{+}\mathbf{v} = \varphi, \quad \mathbf{v}|_{\nu_h} = \nu_0(x).$$

We write the generating scheme for the determination of $y = y(x_i, (j+1)\tau)$, where j+1 is the number of the iteration and τ the iteration parameter, as

$$A^s y_{\overline{t}^s} = \check{\Phi}^s + F^s, \qquad A^s = \prod_{\alpha=1}^p A_{\alpha^s}, \qquad A^s_{\alpha} = E - 0.5 \operatorname{tm}_{s\alpha} \mathring{\Lambda}_{\alpha},$$

$$\check{\Phi}^{s} = \sum_{\alpha=1}^{p} \varkappa_{s\alpha} \mathring{\Lambda}_{\alpha} \check{y}^{s} + (\lambda + \mu) \sum_{k=s+1}^{p} \mathring{\Lambda}_{sk} \check{y}^{k} + \varphi^{s},$$

$$F^{s} = (\lambda + \mu) \sum_{k=1}^{s-1} \mathring{\Lambda}_{sk} y^{k}.$$

In this case the numerical algorithm for alternating directions of Section 6 takes the form

$$A_1^s w_{(1)}^s = \Phi^s + F^s, \quad A_{\alpha}^s w_{(\alpha)}^s = w^s_{(\alpha-1)}, \quad \alpha > 1,$$

 $y^s = y^s + \tau w_{(p)}^s, \quad w_{(\alpha)}^s |_{\gamma_h} = 0,$

i.e. for $w_{(\alpha)}^s$ the boundary conditions are always zero.

For comparison we quote one more numerical algorithm (two-layered)

$$A_1^{s}y_{(i)}^{s} = \tau(\tilde{\Phi^{s}} + F^{s}) + A_{s}\tilde{y^{s}}, \quad A_{\alpha^{s}}y_{(\alpha-1)}^{s} = y_{(\alpha-1)}^{s}, \quad \alpha > 1,$$

$$y_{(\alpha)}^{s} = A_{(\alpha+1)}^{s} \dots A_{p}^{s}v^{s} \quad \text{for} \quad x_{\alpha} = 0, \quad x_{\alpha} = l_{\alpha}, \quad \alpha = 1, 2, \dots, p-1:$$

$$y_{(p)}^{s}|_{V_{h}} = v^{s}.$$

The order of the computation is the same as before: initially the first component $y^{(1)}$ is determined, then the second $y^2(s=2)$ etc.

To find y^s , at all nodes of the network ω_h operations of the order $O(1/h_1h_2\ldots h_p)$ are required. The iteration process converges for $\tau=O(h_*)$, $h_*=\min h_{\alpha}$. The rate of convergence is determined by the number of iterations $v\approx O((1/h_*)\times \ln{(1/\epsilon)})$ for p=2 and $O((1/h_*)^{i_3})\ln{(1/\epsilon)})$ for p=3, where ϵ is the required accuracy.

The evaluation for the number of iterations is obtained by the method of energy inequalities by analogy with [10, 7, 4].

Switching the roles of the operators Λ^- and Λ^+ , we obtain the second iteration scheme II $^+$. The same value is obtained for the rate of convergence of the iterations with it as for the scheme II $^-$ described above. It is hoped that this is the basis which the interchange of these two iteration algorithms II $^-$ and II $^+$ can bring to speeding up to convergence.

REFERENCES

- SAMARSKII, A.A., Local one-dimensional difference schemes for multidimensional hyperbolic equations in an arbitrary region. Zh. vychisl. Mat. mat. Fiz. 4, 4, 638 - 649, 1964.
- SAMARSKII, A.A., Economic difference schemes for parabolic equations with compound derivatives. Zh. vychisl. Mat. mat. Fiz. 4, 4, 753 -759, 1964.
- SAMARSKII, A.A., Economic difference method for the solution of a multi-dimensional parabolic equation in an arbitrary region. Zh. vychisl. Mat. mat. Fiz., 2, 5, 787 - 811, 1962.
- SAMARSKII, A.A., Local one-dimensional schemes on non-uniform networks. Zh. υychisl. Mat. mat. Fiz. 3, 3, 431 - 466, 1963.
- SAMARSKII, A.A., Economic difference schemes for systems of parabolic equations. Zh. vychisl. Mat. mat. Fiz., 4, 5, 927 - 930, 1964.
- 6. KONOVALOV, A.N., Application of resolving methods to the numerical solution of dynamic problems in the theory of elasticity. Zh. vychisl. Mat. mat. Fiz., 4, 4, 760 764, 1964.
- KONOVALOV, A.N., Iteration scheme for the solution of static problems in the theory of elasticity. Zh. vychisl. Mat. mat. Fiz., 4, 5, 942 - 945, 1964.
- 8. SAMARSKII, A.A., Schemes for increasing the order of accuracy for a multi-dimensional equation of heat conduction. Zh. vychisl. Mat. aat. Fiz., 3, 5, 812 840, 1963.
- BEREZIN, I.S. and ZHIDKOV, N.P., Numerical Methods (Metody vychislenii), Vol. 2, Fizmatgiz, Moscow, 1960.
- LEES, M., A note on the convergence of alternating direction methods.
 Math. Comput., 16, 77, 70 75, 1962.